

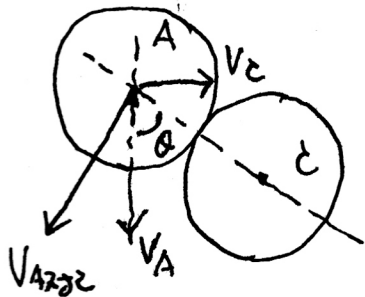
钢丝系:

$$\frac{m\omega^2 x}{mg} = \tan\theta = \frac{dy}{dx}$$

$$\Rightarrow dy = \frac{\omega^2}{g} x dx$$

$$\Rightarrow y = \frac{\omega^2}{2g} x^2$$

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E字理: $mg \cdot 2R (\cos 30^\circ - \cos \theta) = \frac{1}{2} m v_A^2 + m v_C^2$

相对运动: $v_A = v_{A232} \sin \theta$

$$v_C = v_{A232} \cos \theta$$

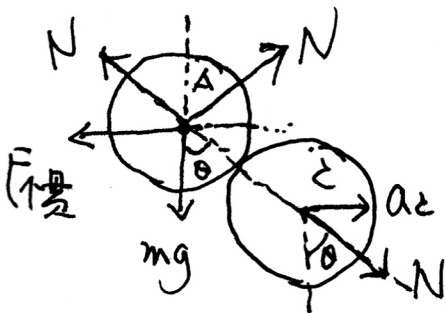
C系中A的受力情况:

$$mg \cos \theta - F_{\text{阻}} \sin \theta - N + N \cos(\pi - 2\theta) = m \frac{v_{A232}}{2R}$$

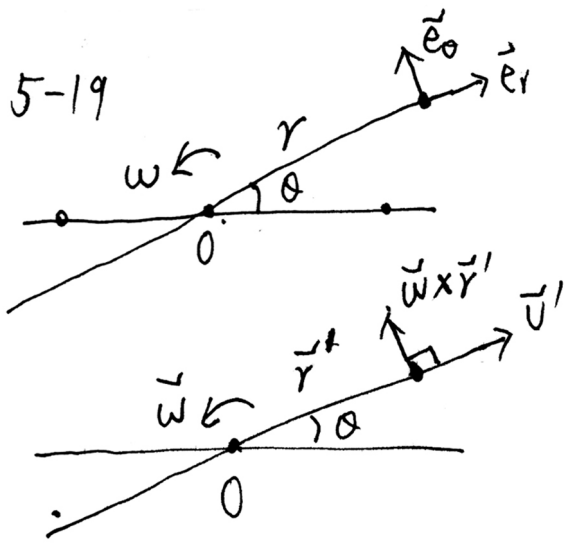
$$F_{\text{阻}} = m a_c$$

$$m a_c = N \sin 2\theta$$

$$\Rightarrow N = \frac{mg \cos \theta}{1 + \sin^2 \theta + \cos^2 \theta} - \frac{2mg (\cos 30^\circ - \cos \theta)}{(1 + \cos^2 \theta)(1 + \sin^2 \theta + \cos^2 \theta)}$$



5-19



运动学:

① 极坐标系: $V_0 = r\dot{\theta}$ $V_r = \dot{r}$

或用② 相对运动关系 (杆绕 O 转动系)

$$\vec{V} = \vec{V}' + \vec{\omega} \times \vec{r}' \quad \vec{r}' = \vec{r}$$

$$V' = \dot{r} = V_r$$

$$|\vec{\omega} \times \vec{r}'| = \omega r = \dot{\theta} r = V_\theta$$

对 O 点 L 守恒:

$$2rmr\dot{\theta} = 2amaw_0 \quad (1)$$

Ek 守恒:

$$2 \frac{1}{2} m (r^2 \dot{\theta}^2 + \dot{r}^2) = 2 \frac{1}{2} m (aw_0)^2 \quad (2)$$

方程①②消去 $\dot{\theta}$ 得:

$$\dot{r} = \frac{aw_0}{r} \sqrt{r^2 - a^2} = \frac{dr}{dt}$$

分离变量, 两边积分:

$$\int_a^r \frac{r dr}{\sqrt{r^2 - a^2}} = \int_0^t aw_0 dt$$

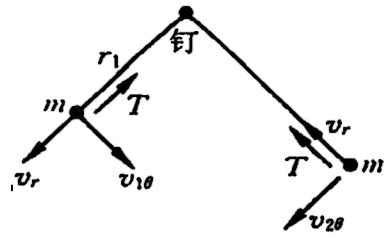
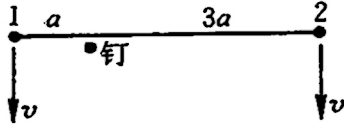
$$\Rightarrow r = \sqrt{a^2 + a^2 w_0^2 t^2}$$

代回①求得:

$$\omega = \dot{\theta} = \frac{w_0}{1 + w_0^2 t^2}$$

$$\alpha = \dot{\omega} = \frac{-2w_0^3 t}{(1 + w_0^2 t^2)^2}$$

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运动学关系 $v_{1r} = v_{2r} = v_r$

各自对钉的角动量守恒

$$mv_{1\theta} r_1 = mva \quad (1)$$

$$mv_{2\theta}(4a - r_1) = mv3a \quad (2)$$

$$\text{动能守恒} \quad \frac{1}{2}m(v_{1\theta}^2 + v_r^2) + \frac{1}{2}m(v_{2\theta}^2 + v_r^2) = 2 \times \frac{1}{2}mv^2 \quad (3)$$

(1) 当球 1 与钉的距离 r_1 达最大时, 必有 $v_r = 0$, 由 (3) 得

$$\frac{1}{2}mv_{1\theta}^2 + \frac{1}{2}mv_{2\theta}^2 = 2 \times \frac{1}{2}mv^2 \quad (4)$$

由 (1) (2) (4) 消去 $v_{1\theta}, v_{2\theta}$, 可得 r_1 的最大值方程

$$\left(\frac{a}{r_1}\right)^2 + \left(\frac{3a}{4a - r_1}\right)^2 = 2$$

(2) 对球 1 有

$$T = -ma_r$$

$$a_r = \frac{d^2 r_1}{dt^2} - r_1 \left(\frac{d\theta}{dt}\right)^2 = \frac{dv_r}{dt} - \frac{v_{1\theta}^2}{r_1} = \frac{dv_r}{dt} - \frac{v^2 a^2}{r_1^3}$$

(1)(2)(3) 联立可得

$$\left(\frac{va}{r_1}\right)^2 + \left(\frac{3va}{4a - r_1}\right)^2 + 2v_r^2 = 2v^2$$

$$\text{两边对 } t \text{ 求导得} \quad \frac{dv_r}{dt} = \frac{v^2 a^2}{2r_1^3} - \frac{9v^2 a^2}{2(4a - r_1)^3}$$

代入 T 表达式得

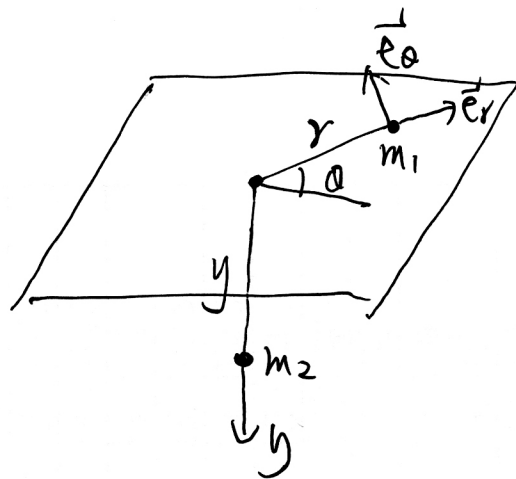
$$T = \frac{mv^2}{2a} \left[\frac{a^3}{r_1^3} + \frac{9a^3}{(4a - r_1)^3} \right]$$

$$T \text{ 极值} \quad T_{\text{Min}} = \frac{(\sqrt{3} + 1)^4}{128} m \frac{v^2}{a}$$

4-25

$$(1) m_2 g = T = m_1 \omega_0^2 r_0 \quad \dots \textcircled{1}$$

$$\omega_0 = \sqrt{\frac{m_2 g}{m_1 r_0}}$$



(2) 几何约束关系:

$$y + r = c$$

$$\Rightarrow a_2 = -\ddot{y} \quad \dots \textcircled{2}$$

牛顿方程:

$$m_2 g - T = m_2 a_2 \quad \dots \textcircled{3}$$

$$-T = m_1 \ddot{r} - m_1 r \dot{\theta}^2 \quad \dots \textcircled{4}$$

对O点, m_1 的 \perp 字恒:

$$m_1 r^2 \dot{\theta} = m_1 r_0^2 \omega_0 \quad \dots \textcircled{5}$$

由①②③④⑤得:

$$(m_1 + m_2) \ddot{y} + m_1 \omega_0^2 r_0 - \frac{m_1 r_0^4 \omega_0^2}{y^3} = 0$$

$$y = r_0 + \delta$$

$$\Rightarrow \ddot{y} = \ddot{\delta}$$

$$(m_1 + m_2) \ddot{\delta} + m_1 \omega_0^2 r_0 - \frac{m_1 r_0^4 \omega_0^2}{(r_0 + \delta)^3} = 0$$

$$(m_1 + m_2) \ddot{\delta} + m_1 \omega_0^2 r_0 \left[1 - \frac{1}{\left(1 + \frac{\delta}{r_0}\right)^3} \right] = 0$$

$$(m_1 + m_2) \ddot{\delta} + m_1 \omega_0^2 r_0 \left[1 - \left(1 - 3\frac{\delta}{r_0}\right) \right] = 0$$

$$\ddot{\delta} + \frac{3m_1}{m_1 + m_2} \omega_0^2 \delta = 0$$

另法:

$$m_2 g = m_1 \omega_0^2 r_0 \quad (1)$$

几何约束:

$$y + r = \zeta_1$$

$$V_2 = -V_1 = -\dot{y} \quad (2)$$

m_1 对 O 点 L 字恒:

$$m_1 r^2 \dot{\theta} = m_1 r_0^2 \omega_0$$

$$\dot{\theta} = \frac{r_0^2 \omega_0}{r^2} \dots (3)$$

$m_1 + m_2$ 的 E 字恒 (以圆周边缘时为势能 0 点)

$$\frac{1}{2} m_2 V_2^2 + \frac{1}{2} m_1 (V_r^2 + V_\theta^2) + m_2 g (r - r_0) = \zeta_2 \quad (4)$$

②③ 代入 ④:

$$\frac{1}{2} (m_1 + m_2) \dot{y}^2 + \frac{1}{2} m_1 r^2 \dot{\theta}^2 + m_1 \omega_0^2 r_0 (r - r_0) = \zeta_2$$

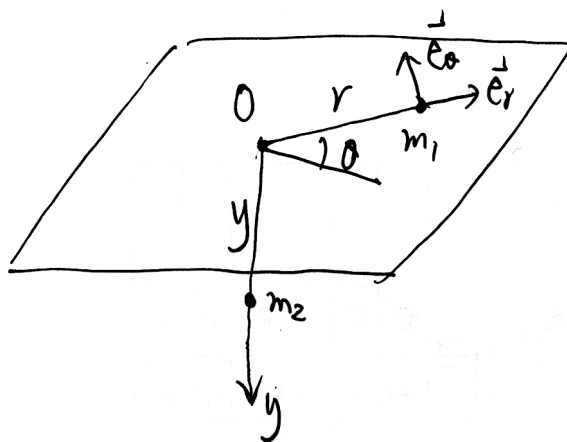
$$\frac{1}{2} (m_1 + m_2) \dot{y}^2 + \frac{1}{2} m_1 \frac{r_0^4 \omega_0^2}{r^2} + m_1 \omega_0^2 r_0 (r - r_0) = \zeta_2 \quad (5)$$

⑤ 式两边对 t 求导:

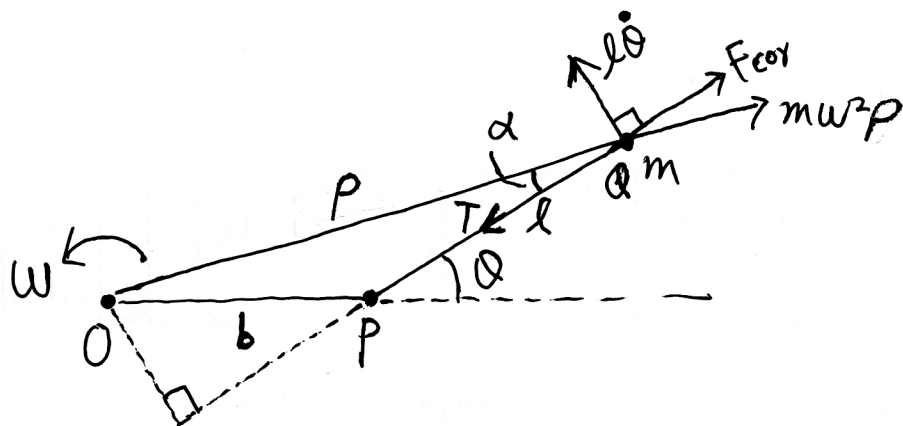
$$(m_1 + m_2) \dot{y} \ddot{y} - m_1 \frac{r_0^4 \omega_0^2}{r^3} \dot{y} + m_1 \omega_0^2 r_0 \dot{y} = 0$$

$$(m_1 + m_2) \ddot{y} + m_1 \omega_0^2 r_0 - \frac{m_1 r_0^4 \omega_0^2}{r^3} = 0$$

后面同前.



4.9



以台面为参考系, 以 P 为参考点, 张力 T 、科氏力 F_{cor} 力矩为 0。
 惯性离心力力矩 (设逆时针为正)

$$M = -l \sin \alpha m \omega^2 \rho = -m \omega^2 l b \sin \theta$$

对 P 点角动量:

$$L = l m l \dot{\theta}$$

由角动量定理 $\frac{dL}{dt} = M$ 得:

$$m l^2 \ddot{\theta} = -m \omega^2 l b \sin \theta$$

$$\text{即 } \ddot{\theta} = -\frac{\omega^2 b}{l} \sin \theta$$

另法: 列极坐标系下横向运动方程:

$$m(2\dot{r}\dot{\theta} + r\ddot{\theta}) = -m\omega^2 \rho \sin \alpha = -m\omega^2 b \sin \theta$$

$$r = l \quad \dot{r} = 0$$

$$\Rightarrow \ddot{\theta} = -\frac{\omega^2 b}{l} \sin \theta$$